



THE UNIVERSITY *of* EDINBURGH  
**informatics**

# Advanced Robotics

## Dynamics I

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# Dynamics vs. Kinematics

- ❑ So far we assumed that we can generate any  $\mathbf{v}_q$  on our robot (eg when looking at forward kinematics)
- ❑ However this is rarely the case, eg:
  - ❑ A flying airplane: You cannot command it to hold still in the air or move straight up
  - ❑ A car: You cannot command it to move sideways
  - ❑ Your arm: You can't command it to throw a ball with arbitrary velocity (force limits)
  - ❑ A torque-controlled robot: You cannot command it to instantaneously change velocity (infinite acceleration/torque)

# Actuation of a robot

An actuator needs a model:

- ❑  $\mathbf{x}$  is the state of the actuator / robot
- ❑  $\mathbf{u}$  is the control input
- ❑ The state at any time depends on both the previous state and the control input:

$$\mathbf{x}_{t+1} = f(\mathbf{x}_t, \mathbf{u}_t) \quad , \text{ with } f \text{ assumed to be smooth}$$

- ❑ Eg: when we looked at forward / inverse kinematics

$$\mathbf{x} = \mathbf{q}, \mathbf{u} = \mathbf{v}_q$$

$$f(\mathbf{q}, \mathbf{v}_q) = \mathbf{q} \oplus \mathbf{v}_q$$

# Three classic models for a robot actuator

□ Velocity source  $\mathbf{x} = \mathbf{q}, \mathbf{u} = \mathbf{v}_q$

Good approximation for hydraulic motors; good for electric actuators only in certain condition (eg industrial manipulators, not legged robots)

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❑ **Acceleration / force source**  $\mathbf{x} = (\mathbf{q}, \mathbf{v}_q), \mathbf{u} = \dot{\mathbf{v}}_q$

Good approximation for electric motors if large contact forces are not involved.

❑ **Torque source**  $\mathbf{x} = (\mathbf{q}, \mathbf{v}_q), \mathbf{u} = \boldsymbol{\tau}$

Good approximation for electric motors. Assumption is that torque is proportional to current. However, gear reductions introduce unmodeled terms that we need to account for.

# Outline

- We discuss the following three topics today:
  - 1D point mass
  - A 'general' dynamic robot ( □ Dynamics II)
  - Joint space control
  
- For now we assume that the robot is fully actuated and that  $\mathbf{v}_q = \dot{\mathbf{q}}$   
(ie velocity and configuration space have the same dimension)
  
- We also assume motors are equipped with accurate **position** sensors (i.e. we know  $\mathbf{q}$  accurately)

# Exercise on the Board

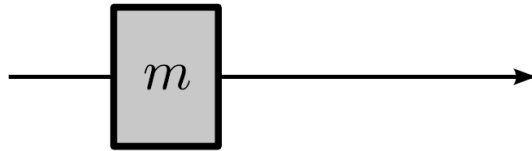
Discuss problem formulation and modelling assumptions associated with moving an object, along a computed trajectory, and controlling against deviations

- How to phrase the questions?
- How does analysis support design?



# Simplest possible case: 1D point mass

- no gravity, no friction



- **State**  $x(t) = (q(t), \dot{q}(t))$  is described by:
  - position  $q(t) \in \mathbb{R}$
  - velocity  $\dot{q}(t) \in \mathbb{R}$
- The **controls**  $u(t)$  is the force we apply on the mass point
- The **system dynamics** is:

$$\ddot{q}(t) = u(t)/m$$

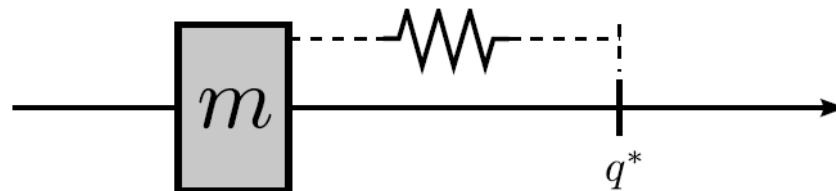
# 1D Mass

- Given current  $q_t$ , what control  $u_t$  to get closer to desired position  $q^*$ ?

# 1D Mass: Proportional Control

- ❑ Given current  $q_t$ , what control  $u_t$  to get closer to desired position  $q^*$ ?
- ❑ Consider an applied force that is an input proportional to the “error”:  
(difference here is that  $u$  is a force and not a velocity)

$$u = K_p (q^* - q)$$



- ❑ You can picture a spring attached to  $q^*$  that pulls the mass towards it. What happens in the absence of friction ?

# 1D Mass: Closed-loop Dynamics

$$m \ddot{q} = u = K_p (q^* - q)$$

$q=q(t)$  is a function of time, this is a second order differential eq.

- Solution: assume  $q(t) = a + b e^{\omega t}$   
(an “non-imaginary” alternative would be  $q(t) = a + b e^{-\lambda t} \cos(\omega t)$ )

$$m b \omega^2 e^{\omega t} = K_p q^* - K_p a - K_p b e^{\omega t}$$

$$(m b \omega^2 + K_p b) e^{\omega t} = K_p (q^* - a)$$

$$\Rightarrow (m b \omega^2 + K_p b) = 0 \wedge (q^* - a) = 0$$

$$\Rightarrow \omega = i \sqrt{K_p/m}$$

$$q(t) = q^* + b e^{i \sqrt{K_p/m} t}$$

This is an oscillation around  $q^*$  with amplitude  $b = q(0) - q^*$  and frequency  $\sqrt{K_p/m}$ !

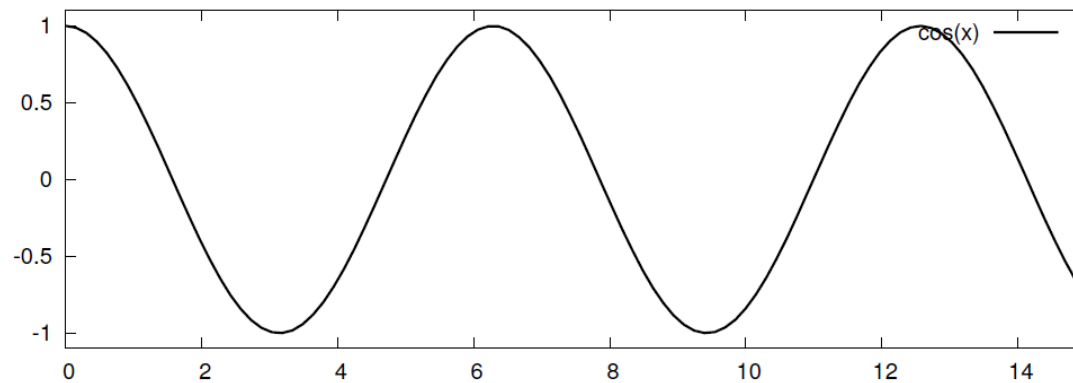
# 1D Mass: Closed-loop Dynamics

What's the effect?

$$m \ddot{q} = u = K_p (q^* - q)$$

$$q(t) = q^* + b e^{i\sqrt{K_p/m} t}$$

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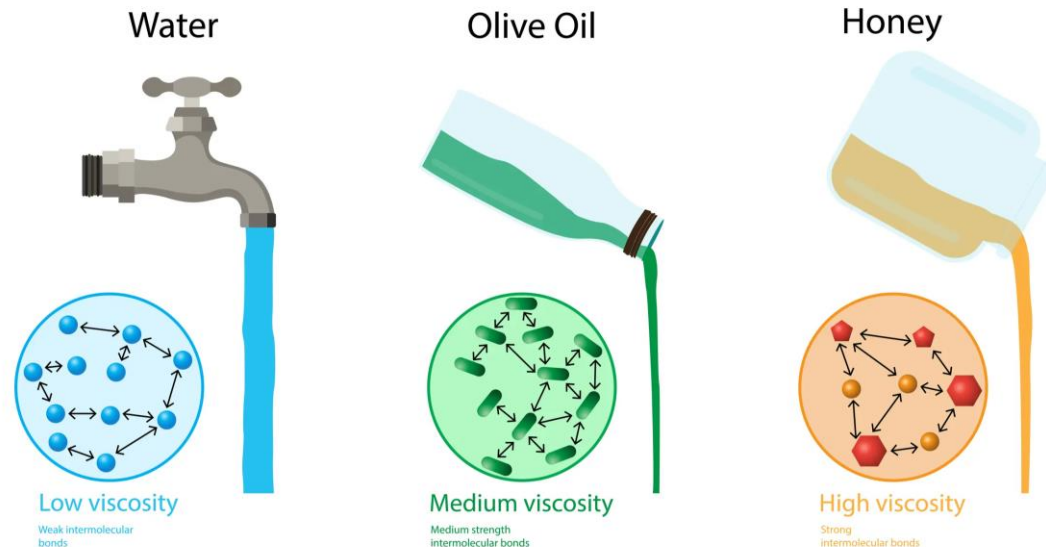


# 1D Mass: Damping forces

Can we shape the dynamics further?

*“Pull less, when we’re heading the right direction already:”*

*“Damp the system:”*



# 1D Mass: Damping forces

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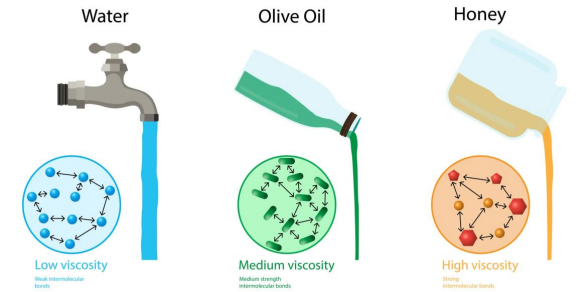
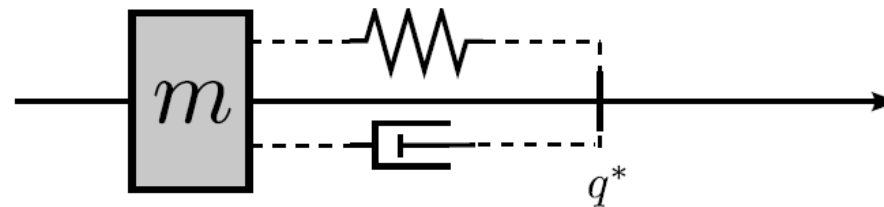
*“Pull less, when we’re heading the right direction already:”*

*“Damp the system:”*

$$u = K_p(q^* - q) + K_d(\dot{q}^* - \dot{q})$$

$\dot{q}^*$  is a desired goal velocity

For simplicity we set  $\dot{q}^* = 0$  in the following.



# 1D Mass: Damping forces

What is the effect?  $m\ddot{q} = u = K_p(q^* - q) + K_d(0 - \dot{q})$

- Solution: again assume  $q(t) = a + be^{\omega t}$

$$m b \omega^2 e^{\omega t} = K_p q^* - K_p a - K_p b e^{\omega t} - K_d b \omega e^{\omega t}$$

$$(m b \omega^2 + K_d b \omega + K_p b) e^{\omega t} = K_p (q^* - a)$$

$$\Rightarrow (m \omega^2 + K_d \omega + K_p) = 0 \wedge (q^* - a) = 0$$

$$\Rightarrow \omega = \frac{-K_d \pm \sqrt{K_d^2 - 4mK_p}}{2m}$$

$$q(t) = q^* + b e^{\omega t}$$

The term  $-\frac{K_d}{2m}$  in  $\omega$  is real  $\leftrightarrow$  exponential decay (damping)



# 1D Mass: Damping forces

What's the effect?

$$q(t) = q^* + b e^{\omega t}, \quad \omega = \frac{-K_d \pm \sqrt{K_d^2 - 4mK_p}}{2m}$$

- Effect of the second term  $\sqrt{K_d^2 - 4mK_p}/2m$  in  $\omega$ :

$$K_d^2 < 4mK_p \Rightarrow \omega \text{ has imaginary part}$$

oscillating with frequency  $\sqrt{K_p/m - K_d^2/4m^2}$

$$q(t) = q^* + b e^{-K_d/2m t} e^{i\sqrt{K_p/m - K_d^2/4m^2} t}$$

$$K_d^2 > 4mK_p \Rightarrow \omega \text{ real}$$

strongly damped

$$K_d^2 = 4mK_p \Rightarrow \text{second term zero}$$

only exponential decay

# 1D Mass: Concept of *damping ratio*

## Alternative Parameterisation

Instead of specifying the *coefficients*  $K_p$  and  $K_d$

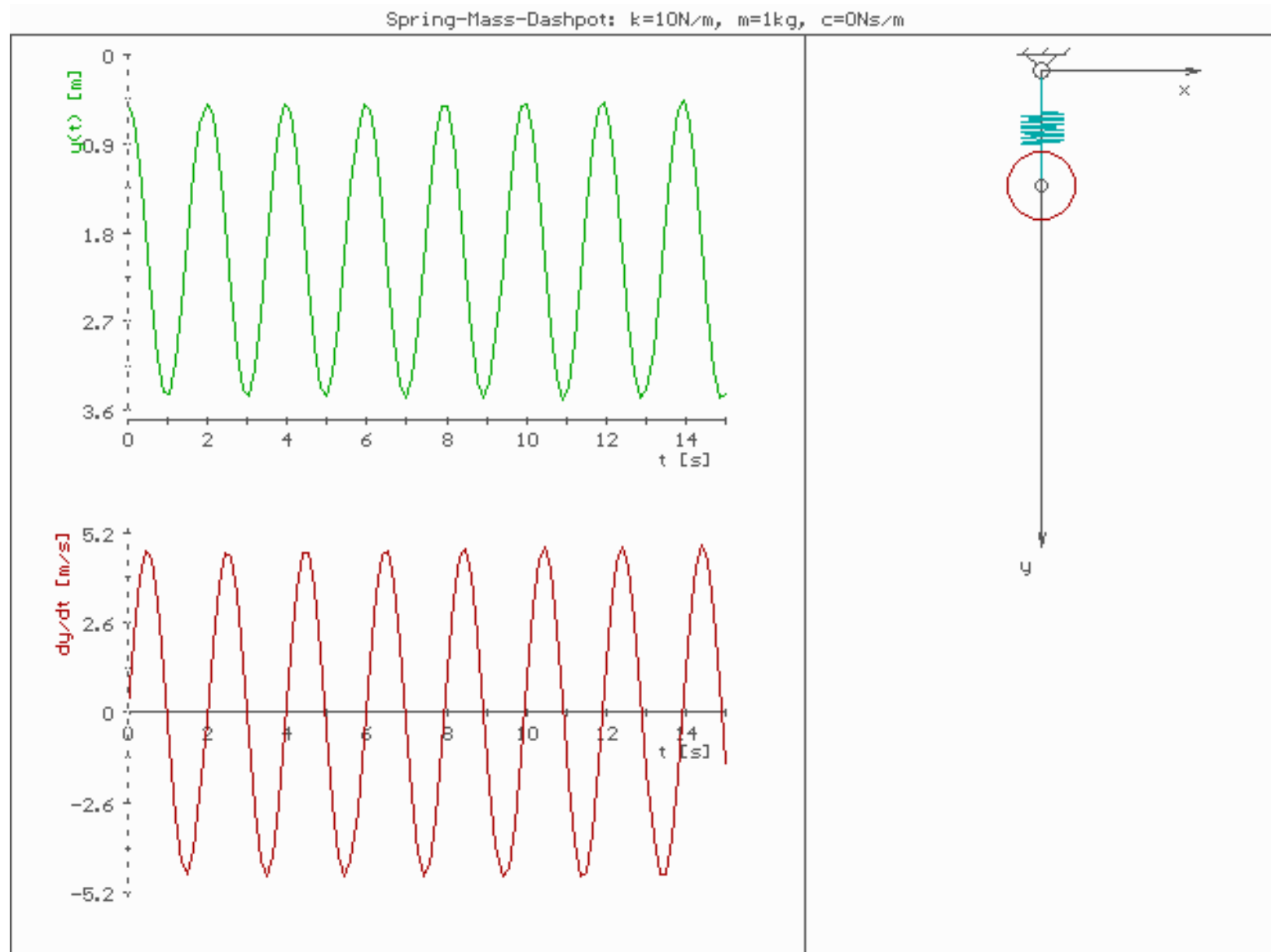
- “wave length”  $\lambda = \frac{1}{\omega_0} = \frac{1}{\sqrt{K_p/m}}$ ,  $K_p = m/\lambda^2$
- damping ratio  $\xi = \frac{K_d}{\sqrt{4mK_p}} = \frac{\lambda K_d}{2m}$ ,  $K_d = 2m\xi/\lambda$

$\xi > 1$ : over-damped

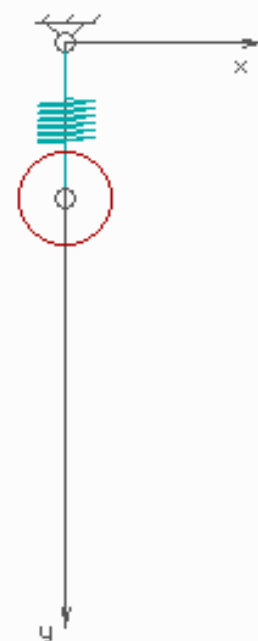
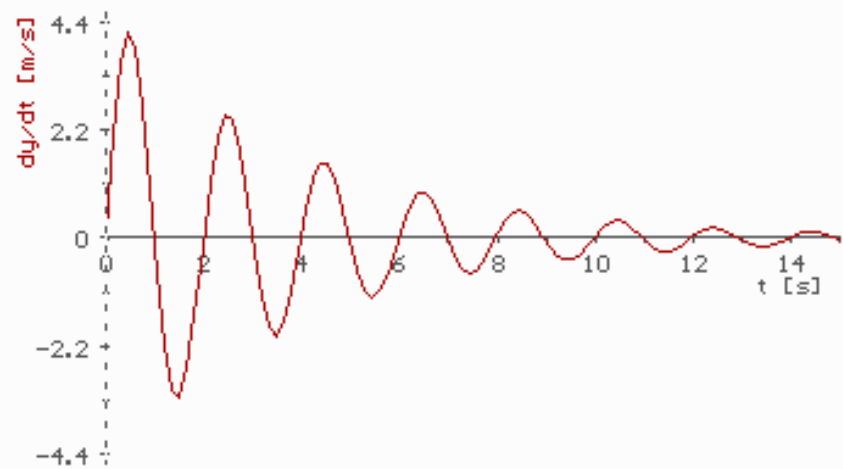
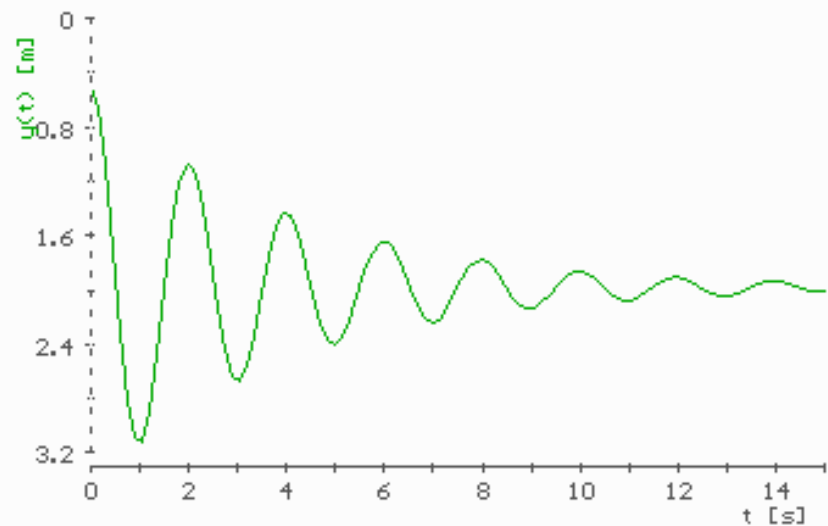
$\xi = 1$ : critically damped

$\xi < 1$ : oscillatory-damped

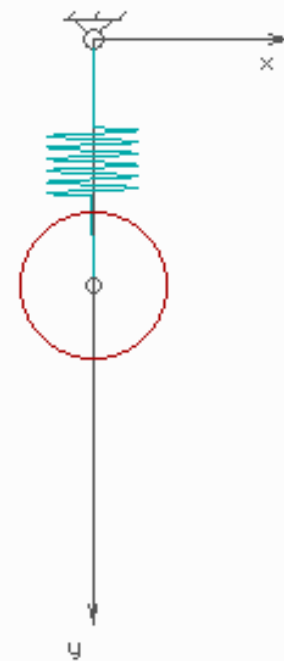
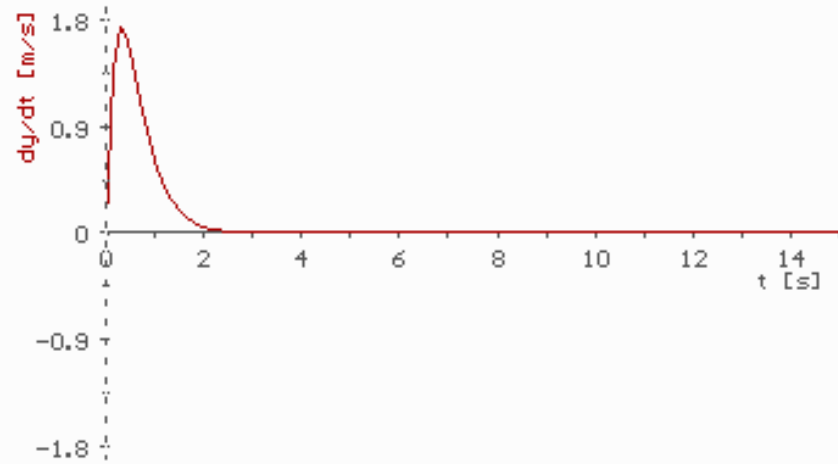
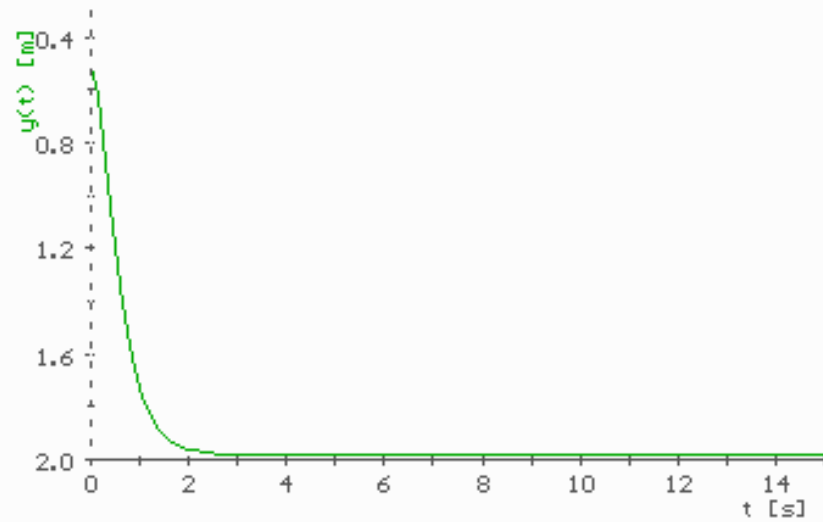
$$q(t) = q^* + be^{-\xi/\lambda t} e^{i\omega_0\sqrt{1-\xi^2} t}$$



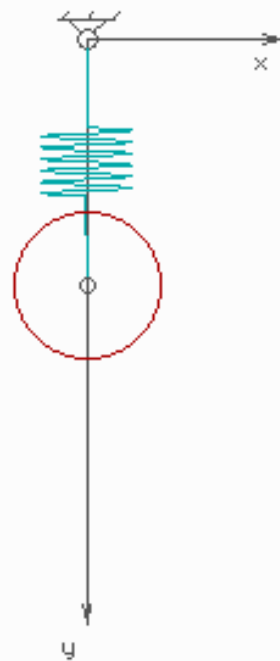
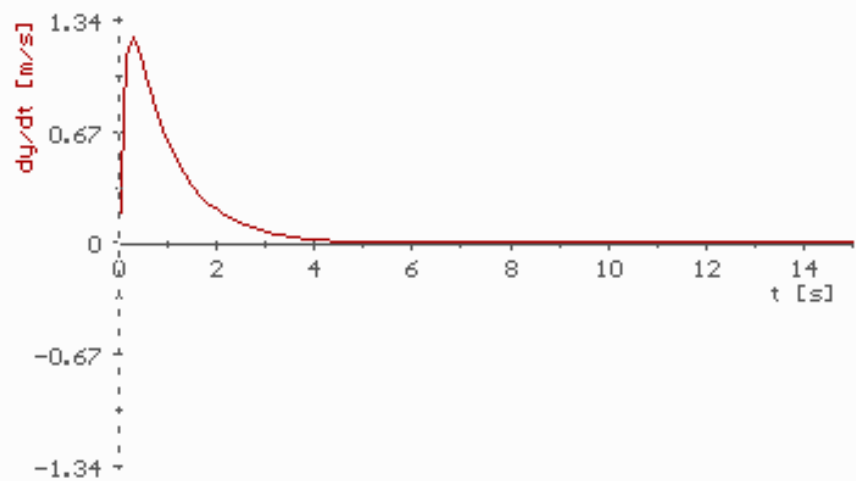
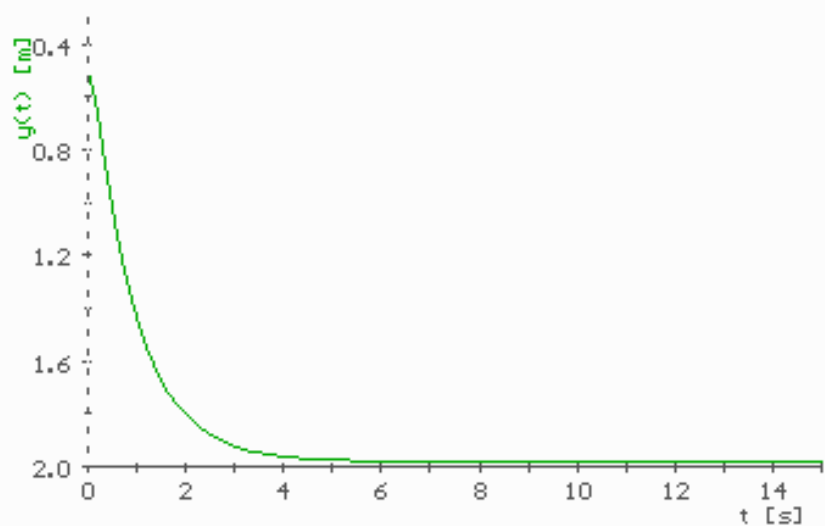
Spring-Mass-Dashpot:  $k=10\text{N/m}$ ,  $m=1\text{kg}$ ,  $c=0.5\text{Ns/m}$



Spring-Mass-Dashpot:  $k=10\text{N/m}$ ,  $m=1\text{kg}$ ,  $c=6.3245\text{Ns/m}$



Spring-Mass-Dashpot:  $k=10\text{N/m}$ ,  $m=1\text{kg}$ ,  $c=10\text{Ns/m}$



# Adding another term to this Spring-damper system

$$u = K_p(q^* - q) + K_d(\dot{q}^* - \dot{q}) + K_i \int_{s=0}^t (q^*(s) - q(s)) ds$$

- **PID control**

- Proportional Control (“Position Control”)

$$f \propto K_p(q^* - q)$$

- Derivative Control (“Damping”)

$$f \propto K_d(\dot{q}^* - \dot{q}) \quad (\dot{x}^* = 0 \rightarrow \text{damping})$$

- Integral Control (“Steady State Error”)

$$f \propto K_i \int_{s=0}^t (q^*(s) - q(s)) ds$$

# 1D Mass: Summary

- ❑ Dynamics of a 1D mass-spring-damper system, a spring and a damper added to a point mass, i.e. spring and damping forces (aka PD controller)
- ❑ Resultant force acting on the system in a linear ‘control law’

$$\pi : (q, \dot{q}) \mapsto u = K_p(q^* - q) + K_d(\dot{q}^* - \dot{q})$$

(linear in the *dynamic system state*  $x = (q, \dot{q})$ )

- ❑ With such simple linear rules, we can modulate the dynamic response of the system by tuning the ‘strength’ of spring and damper (ie, PD gains in the PD ‘control law’)
- ❑ \*trade-off: there is no optimality criterion supporting such rules and the resulting motions (hence, we may be able to do better)