# Quantum Cyber Security Lecture 8: Quantum Key Distribution II

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- From **QBER** to secure key distribution (general expression and how to use it)
- **Simplifying assumptions** (physical restrictions, classical efficiency, adversary's limitations, composability)
- Security proof for the basic BB84 protocol
- Classical post-processing and its cost

• General Expression:

$$R = \frac{Q}{2} \left( \xi H(A:B) - S(A:E) - \Delta(n,\epsilon) \right)$$

• *R* is the secret key-rate: Expected secret bits per qubit sent.

#### • General Expression:

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- *R* is the secret key-rate: Expected secret bits per qubit sent.
- *Q* is the prob that sent single-photons are detected (not lost)
- factor <sup>1</sup>/<sub>2</sub> is due to the raw key that includes only the positions that Alice and Bob measured in same basis
- $\xi$  is due to non-ideal classical post-processing (IR and PA)
- Δ(n, ε) is a factor due to finite-size effects (measured value differing from expectation)

• For simplicity we consider: perfect detection, ideal post-processing and asymptotic limit

Q = 1;  $\xi = 1$ ;  $\Delta(n, \epsilon) = 0$ 

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• Using details of BB84 protocol we get:

$$R_{\rm BB84} = \frac{1}{2} \left( 1 - h(e_b) - h(e_p) \right)$$
 (1)

where  $e_b$  and  $e_p$  are the average errors in the  $\{|0\rangle, |1\rangle\}$  and  $\{|+\rangle, |-\rangle\}$  bases and  $h(p) := -p \log_2 p - (1-p) \log_2(1-p)$  is the binary entropy

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• If the errors in different bases equal and equal to the QBER  $(e_b = e_p = D)$  we finally get:

$$R_{\rm BB84} = \frac{1}{2} \left( 1 - 2h(D) \right)$$
 (2)

#### Examples: How to compute key rate (ideal case)

• Example 1: Given  $e_b = 0.05$ ,  $e_p = 0.1$  find the rate.

$$R_{\rm BB84} = \frac{1}{2} \left( 1 - h(0.05) - h(0.1) \right) = \frac{1}{2} \left( 1 - 0.29 - 0.47 \right) = 0.12$$

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• Example 2: Which is the largest QBER *D* (symmetric in two bases) that BB84 does not abort (error tolerance)?

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• Example 3: Does intercept, measure Z & resend attack abort?

$$e_b = 0$$
;  $e_p = 0.5$ ;  $R_{BB84} = \frac{1}{2}(1 - h(0.5)) = 0$ 

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  - Fully trusted quantum devices: Assumptions on how the preparation and measuring devices behave and what information on their workings could leak (e.g. due to a hacking/side-channel attack)

• Finite-size effects: Bounds on the mutual information are computed based on expectations values of observables.

Measured values differ from expectation values for finite size keys, but they converge (exponentially – cf Chernoff bounds) when the length of the string tends to infinity. Finite-size corrections are required for practical QKD

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• Cost of classical post-processing: Theoretical error-correction (IR) leaks information to make A', B' perfectly correlated, related with the conditional entropy H(A|B)

**Practical error-correction** leaks more bits of information (cf  $\xi$ -coefficient)

- Ability of adversary: (from weaker to stronger)
  - i.i.d. attacks: Interacts with sent each qubit separately, independently and identically Can reduce remarks regarding strings of qubits to the expected effect on a single qubit

State Alice prepares:  $|x\rangle_A \langle x| \otimes \rho_B^x$  where x represents the classical info Alice stores (which BB84 state was prepared).

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 σ<sup>x</sup><sub>E</sub> = Tr<sub>B</sub>(σ<sup>x</sup><sub>BE</sub>) is Eve's system. She performs measurement to obtain the max info on classical variable x

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- Collective attacks: Uses different private system for each qubit, interacts with each qubit (non iid) and then measures conditionally on other previous actions
- Coherent Attacks: Uses private system(s), interacts with all passing qubits, stores everything and measures all systems at the end (possibly in entangled basis)

 Composability: In modern crypto, security is proven is such a way that essential properties proven are directly maintained when composed with other protocols (e.g. used as subroutine in a larger protocol)

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- Way to prove:
  - Define ideal properties that protocol would have
  - Any adversary has bound probability of distinguishing the real protocol from a simulated protocol that uses the ideal protocol
  - In quantum case, bounding this probability reduces in bounding the **trace-distance** of the real protocol from an ideal protocol

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- Simplifying Assumptions:
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  - No losses (Q = 1)
  - trusted and ideal single-photon source and measuring devices
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- Adversarial Model: i.i.d. and non-composable
- Proof **can be generalised** for stronger adversaries and without the simplifying assumptions, adjusting parameters and with simple protocol modifications

- i.i.d. case see effects on single qubit (rather than strings)
- Need to bound (subject to average errors  $e_b, e_p$ ):

$$R = \frac{1}{2}(H(A:B) - S(A:E))$$

• See also alternative proof later

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- See also alternative proof later
- $H(A:B) = H(A) H(A|B) = 1 \frac{1}{2}(h(e_b) + h(e_p))$  H(A) = 1 since A is chosen randomly  $H(A) = -1/2 \log_2 \frac{1}{2} - 1/2 \log_2 \frac{1}{2} = 1$  H(A|B) when state is sent in the Z basis is  $H(A|B) = -(1 - e_b) \log_2(1 - e_b) - e_b \log_2 e_b = h(e_b)$  and happens in half cases H(A|B) when state is sent in the X basis is  $H(A|B) = -(1 - e_p) \log_2(1 - e_p) - e_p \log_2 e_p = h(e_p)$  and happens in the other cases
- Overall:  $H(A|B) = \frac{1}{2}(h(e_b) + h(e_p))$

• Need to bound S(A : E). Eve has the quantum state:

 $\sigma_{E}^{\mathsf{x}} = \operatorname{Tr}_{B}\left(U_{BE}(\rho_{B}^{\mathsf{x}} \otimes |0,0\rangle_{E}\langle 0,0|\right)\right)$ 

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Need to min the classical info about x that she can extract. Accessible Information: Given ensemble  $F := \{(p(x), \sigma^x)\}$ , the (generalised) measurement  $\{M\}$ , and the random variable corresponding to the measurement's outcome  $Y_M$ :

 $I_{acc}(F) = \max_{M} H(X : Y_{M})$ 

Holevo bound: Given ensemble F := {(p(x), σ<sup>x</sup>)}, the accessible information is upper bounded by the Holevo quantity χ(F)

$$I_{acc}(F) \le \chi(F) := S(\sum_{x} p(x)\sigma^{x}) - \sum_{x} p(x)S(\sigma^{x})$$

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- In our case Eve's register is  $d = 2^2$  and thus:

#### $I_{acc} \leq \chi(F) \leq 2$

The maximum classical information extractable from a single qubit (irrespective of the number of classical states encoded) is one bit!

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• Let  $\sigma_{BE} = |\psi\rangle_{BE} \langle \psi|$  be a pure state (global), then:  $S(\sigma_B) := S(\operatorname{Tr}_E(\sigma_{BE})) = S(\sigma_E) := S(\operatorname{Tr}_B(\sigma_{BE}))$ 

(See Schmidt decomposition for proof)

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In our case (F = {p(x), σ<sub>E</sub><sup>x</sup>}) for the individual terms, σ<sub>BE</sub><sup>x</sup> is pure so we can use the entropy of the B system, which for given x is give by the resp error:

$$S(A:E) \leq I_{acc}(F) \leq \chi(F) = S(\sigma_E) - \frac{1}{2}(h(e_b) + h(e_p))$$

Leading to

$$R \geq \frac{1}{2} \left( H(A:B) - \chi(F) \right) = \frac{1}{2} \left( 1 - S(\sigma_E) \right)$$

# BB84: A basic security proof

- It can also be shown that  $S(\sum_x \frac{1}{4}\sigma_E^x) \le h(e_b) + h(e_p)$
- Algebraically has 2 maximum value (if Bob's state is random and independent of the state of Alice)
- This leads to the final expression given by Eq. (1)

$$R_{BB84} \geq \frac{1}{2} \left(1 - h(e_b) - h(e_p)\right)$$

- This becomes negative if  $e_p$ ,  $e_b$  increase (has max value -1 when these become 1/2)
- The overall 1/2 factor in Eq. (1) can be removed if the states sent are mainly in one (preferred) basis. This is possible if there are sufficient states sent in the other basis to have good enough statistics (cf finite-size effects)

- Consider tripartite system ABE. System A is either measured in Z or X basis to result to classical variable A<sup>Z</sup>, A<sup>X</sup> resp
- For simplicity systems A, B are assumed to be single qubits, then the following inequality holds for all global states ρ<sub>ABE</sub>

$$S(A^X|B) + S(A^Z|E) \ge 1$$
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We have

H(A:B) - S(A:E) = S(A|E) - H(A|B) = S(A|E) - S(A|B)

which we can break to two terms depending the basis used:

$$\frac{1}{2}\left(S(A^{Z}|E) - S(A^{Z}|B^{Z}) + S(A^{X}|E) - S(A^{X}|B^{X})\right)$$

• From Eq. (3) we get:

$$I(A:B) - S(A:E) \ge \frac{1}{2} \left( (1 - S(A^X|B^X)) - S(A^Z|B^Z) + (1 - S(A^Z|B^Z)) - S(A^X|B^X) \right)$$

• From Eq. (3) we get:

 $I(A:B) - S(A:E) \ge \frac{1}{2} \left( (1 - S(A^X | B^X)) - S(A^Z | B^Z) + (1 - S(A^Z | B^Z)) - S(A^X | B^X) \right)$ • Noting that  $S(A^Z | B^Z) = h(e_b)$ ;  $S(A^X | B^X) = h(e_p)$  $I(A:B) - S(A:E) \ge 1 - h(e_b) - h(e_p)$ 

which then leads to the known expression Eq. (1)

## Classical Post-Processing

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Privacy Amplification (PA): Use family of universal hash functions to ensure that the final (smaller) key Alice and Bob share, is completely secret from Eve (i.e. amplify the privacy). Map strings to smaller strings s.t. entropy H(A''|E'') of new strings A'' = g(A'); E'' = g(E') is maximum

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Due to non-ideal procedure, to ensure identical output leaked bits are increased by a factor  $\xi$  compared to ideal Shannon limit

• Leftover hash lemma: if a secret bit-string A of length n has t bits leaked (at unknown positions), then you can produce a bit string of  $m \le n - t - 2\log_2(1/\epsilon)$  bits that is totally secret (almost optimal)

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- 2-Universal hash family: Let a family of functions  $g_i \in G$  with  $i \in S$  (cardinality of family |S|), where  $g_i : \{U \to [m] = \{0,1\}^m\}$ :
  - for fixed  $A \in U$  if  $g_i$  is randomly chosen from the family, the  $g_i(A)$  is uniformly distributed in [m]
  - ② for any pair A, E ∈ U, if i is chosen randomly,  $g_i(A), g_i(E)$  are independent variables

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  - ② for any pair A, E ∈ U, if i is chosen randomly,  $g_i(A), g_i(E)$  are independent variables
- Consider a string A with (n t)-bits of randomness. If

 $m \leq (n-t)$  then using the 2-universal hash family G:

#### $\delta[(g_i(A),i),(R,i)] \leq \epsilon$

*R* uniformly random *m*-bit string,  $\delta$  stat distance