Quantum Cyber Security Lecture 4: Quantum Information Basics II

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- Understand mathematics of quantum states
 - Most general way to describe quantum systems

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- Quantum measurements and their mathematics

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 Most general way to describe quantum systems
- Quantum measurements and their mathematics
- Quantum operations and their mathematics
- Properties and concepts of classical and quantum information theory



Describe



Observe



Evolve



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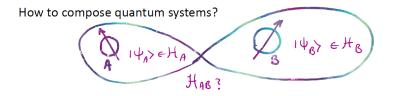
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- Or multiple qubits: n-qubits have states of dimension d = 2ⁿ
 Hilbert space
- Let's see how to compose quantum systems (e.g. two qubits)



 Two Hilbert spaces *H_A*, *H_B* can form a new (composite) Hilbert space *H_{AB}*

 $\dim \mathcal{H}_{AB} = \dim \mathcal{H}_A \times \mathcal{H}_B$

- Basis vectors of composite are the "product" of the basis vectors of the individual spaces
- Tensor product $\mathcal{H}_{AB} := \mathcal{H}_A \otimes \mathcal{H}_B$

Let V and W be two vector spaces with dim m and n. The tensor product $V \otimes W$ of these vector spaces is a vector space of dimension $m \times n$ to which is associated a bilinear map that maps a pair $(v, w), v \in V, w \in W$ to an element of $V \otimes W$ denoted as $v \otimes w$.

- Let $|i\rangle$ and $|j\rangle$ be orthonormal bases for V and W respectively Then $|i\rangle \otimes |j\rangle$ is orthonormal basis for $V \otimes W$
- General state $|\psi\rangle_{VW} = \sum_{i,j} \psi_{ij} |i\rangle_V \otimes |j\rangle_W$

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- Matrix rep. of operator tensor products:

Let A_{ij} matrix elements of A and B_{kl} of B: $A \otimes B = \sum i, j, k, lA_{ij}B_{kl} |i\rangle \langle j| \otimes |k\rangle \langle l|$

$$A \otimes B = \begin{bmatrix} A_{11}B & A_{12}B & \cdots & A_{1n}B \\ A_{21}B & A_{22}B & \cdots & A_{2n}B \\ \vdots & \vdots & \vdots & \vdots \\ A_{m1}B & A_{m2}B & \cdots & A_{mn}B \end{bmatrix}$$

• Dirac notation:

 $|0
angle\otimes|+
angle,\,|angle\otimes|angle\otimes|+
angle,\,|01
angle\otimes|angle$

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Matrix notation:

$$|0\rangle \otimes |0\rangle = |00\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \otimes \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \times \begin{pmatrix} 1 \\ 0 \\ 0 \times \begin{pmatrix} 1 \\ 0 \end{pmatrix} \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}$$

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• Operators:

$$\begin{pmatrix} 1 & 2 \\ 0 & 4 \end{pmatrix} \otimes \begin{pmatrix} 2 & 1 \\ 3 & i \end{pmatrix}$$

Properties of Tensor Products

- Properties:
 - $c(|v\rangle \otimes |w\rangle) = (c|v\rangle) \otimes |w\rangle = |v\rangle \otimes (c|w\rangle)$ where c is a scalar.
 - $(|v_1\rangle + |v_2\rangle) \otimes |w\rangle = |v_1\rangle \otimes |w\rangle + |v_2\rangle \otimes |w\rangle$
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- Tensor product isn't commutative |v⟩ ⊗ |w⟩ ≠ |w⟩ ⊗ |v⟩
 (not the order of the spaces is conventional, could reorder them if needed, but on all terms of one expression!)
- A vector tensored k-times with itself: $|\psi\rangle^{\otimes k}$

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- If A acts on V and B acts on W, then

 $(A \otimes B)(|v\rangle \otimes |w\rangle) = A |v\rangle \otimes B |w\rangle$

Examples: tensor product operators

•
$$O_1 |\psi_1\rangle = |\phi_1\rangle$$
 and $O_2 |\psi_2\rangle = |\phi_2\rangle$
 $(O_1 \otimes O_2)(|\psi_1\rangle \otimes |\psi_2\rangle) = O_1 |\psi_1\rangle \otimes O_2 |\psi_2\rangle = |\phi_1\rangle \otimes |\phi_2\rangle$

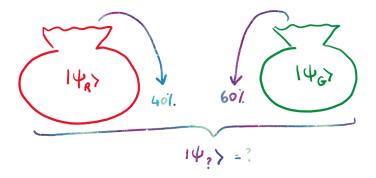
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• $X |j\rangle = |j \oplus 1\rangle$, $Z |j\rangle = (-1)^j |j\rangle$
 $(X \otimes Z) |01\rangle = X |0\rangle \otimes Z |1\rangle = |1\rangle \otimes (-1) |1\rangle = -|11\rangle$

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• $H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$, $X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$
 $H \otimes X = \frac{1}{\sqrt{2}} \begin{pmatrix} X & X \\ X & -X \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 & 1 \\ 1 & 0 & 1 & 0 \\ 0 & 1 & 0 & -1 \\ 1 & 0 & -1 & 0 \end{pmatrix}$

Ensembles of Quantum States



• Examples:

$$egin{aligned} &
ho = rac{1}{2} \left| 0
ight
angle \left\langle 0
ight| + rac{1}{2} \left| +
ight
angle \left\langle +
ight| \ &
ho = p_1 \left| \psi_R
ight
angle \left\langle \psi_R
ight| + p_2 \left| \psi_G
ight
angle \left\langle \psi_G
ight| \ &
m where \ p_1 = 0.4 \ , \ p_2 = 0.6 \end{aligned}$$

Density Matrices (Recall)

Definition: A density matrix is a matrix (or operator) ρ that: 1 is Hermitian $\rho^{\dagger} = \rho$

2 positive semi-definite (i.e. has non-negative eigenvalues)

3 has unit trace $Tr(\rho) = 1$

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- Real eigenvalues, non-negative, normalised
- Pure state vector $|\psi\rangle$ goes to pure density matrix: $\rho_{\psi}:=|\psi\rangle\left\langle\psi\right|$
- Can incorporate probabilities over quantum states (ensembles)

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Express the composite $ho_A \otimes
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• All mixed states can be expressed as ensembles (diagonalise!)

Entanglement

• There are states that CANNOT be written as tensor product of individual states

Product states: $|\psi_{AB}\rangle = |\phi_A\rangle \otimes |\phi_B\rangle$

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- Mixed states that cannot be written as mixtures of product states are **entangled**

Mixture of product states: $\rho = \sum_{ij} p_{ij} \rho_i \otimes \rho_j$

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• To describe a subsystem of a (pure or mixed) entangled state, we need density matrices!

• $|\Phi^+
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ight)$ Cannot be written as product state

Bell states are entangled

• $|\Phi^+\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$ Cannot be written as product state Proof: Assume that there exists states $|a\rangle = a_0 |0\rangle + a_1 |1\rangle$ and $|b\rangle = b_0 |0\rangle + b_1 |1\rangle$ such that $|\Phi^+\rangle = |a\rangle \otimes |b\rangle$ • $|\Phi^+\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$ Cannot be written as product state Proof: Assume that there exists states $|a\rangle = a_0 |0\rangle + a_1 |1\rangle$ and $|b\rangle = b_0 |0\rangle + b_1 |1\rangle$ such that $|\Phi^+\rangle = |a\rangle \otimes |b\rangle$

It follows that

 $\ket{a}\otimes\ket{b}=a_{0}b_{0}\ket{00}+a_{0}b_{1}\ket{01}+a_{1}b_{0}\ket{10}+a_{1}b_{1}\ket{11}$

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Since there are no $|01\rangle$, $|10\rangle$ terms, we know that $a_0b_1 = a_1b_0 = 0$

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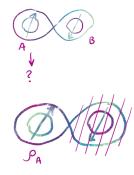
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Since there are no $\ket{01}, \ket{10}$ terms, we know that $a_0b_1 = a_1b_0 = 0$

(i) if $a_0 = 0$ then the term that has $|00\rangle$ vanishes (which it shouldn't)

(ii) if $b_1 = 0$ the term with $|11\rangle$ vanishes \Box

Subsystems and partial trace



• Let
$$\rho_{AB} = |\psi_{AB}\rangle \langle \psi_{AB}|$$

Then $\rho_A := \operatorname{Tr}_B(\rho_{AB})$ and $\rho_B := \operatorname{Tr}_A(\rho_{AB})$

• It easy to see that for product states $\rho_{AB} = \rho_A \otimes \rho_B$ this is the case

Partial Trace

- Consider $M_{AB} = \sum_{i,j,k,l} c_{ijkl} |i\rangle \langle j|_A \otimes |k\rangle \langle l|_B$
- Partial trace over B:

$$M_{A} := \operatorname{Tr}_{B}(M_{AB}) = \sum_{i,j,k,l} c_{ijkl} |i\rangle \langle j|_{A} \times \operatorname{Tr}(|k\rangle \langle l|_{B})$$

(note the trace is number not a matrix)
$$= \sum_{i,j,k,l} c_{ijkl} \langle i|_{A} \times (\langle l|k\rangle_{B})$$
 (using cyclic property)

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$$= \sum_{i,j,k,l} c_{ijkl} |i\rangle \langle j|_A \times (\langle l|k\rangle_B) \text{ (using orthogonality)}$$
$$\sum_{i,j} \sum_k c_{ijkk} |i\rangle \langle j|_A$$

Partial Trace

- Consider $M_{AB} = \sum_{i,j,k,l} c_{ijkl} \ket{i} \bra{j}_A \otimes \ket{k} \bra{l}_B$
- Partial trace over B:

$$\begin{split} M_A &:= \operatorname{Tr}_B(M_{AB}) = \sum_{i,j,k,l} c_{ijkl} |i\rangle \langle j|_A \times \operatorname{Tr}(|k\rangle \langle l|_B) \\ \text{(note the trace is number not a matrix)} \end{split}$$

- $= \sum_{i,j,k,l} c_{ijkl} |i\rangle \langle j|_A \times (\langle l|k\rangle_B) \text{ (using cyclic property)}$ = $\sum_{i,j,k,l} c_{ijkl} |i\rangle \langle j|_A \times (\langle l|k\rangle_B) \text{ (using orthogonality)}$ $\sum_{i,j} \sum_k c_{ijkk} |i\rangle \langle j|_A$
- Reduced matrix (partial trace over B) is a matrix at space A
- Partial trace over A is defined similarly

• Reduced state for product states $\rho_{AB} = \rho_A \otimes \rho_B$ As expected: ρ_A and ρ_B

- Reduced state for product states $\rho_{AB} = \rho_A \otimes \rho_B$ As expected: ρ_A and ρ_B
- Reduced state for entangled (Bell) state:

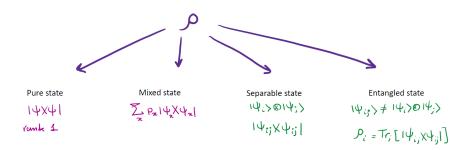
$$\rho_{AB} = |\Phi^+\rangle \langle \Phi^+|$$

$$|\Phi^+\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$

$$\rho_A = \operatorname{Tr}(\rho_{AB}) = \dots = \frac{1}{2} (|0\rangle \langle 0| + |1\rangle \langle 1|)_A$$

$$\rho_B = \operatorname{Tr}(\rho_{AB}) = \dots = \frac{1}{2} (|0\rangle \langle 0| + |1\rangle \langle 1|)_B$$

One density operator to rule them all!



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- It generalises for observable O

Born Rule:

The measured result for an observable O, on a quantum system $|\psi\rangle$ is given by its eigenvalues λ The probability of getting a specific eigenvalue λ_i is equal to $\mathbf{p}(i) = \langle \psi | P_i | \psi \rangle$ or more generally for a density matrix ρ is given by $p(i) = Tr[P_i\rho P_i^{\dagger}]$ Where P_i is the projection onto the eigenspace of O corresponding to λ_i

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or more generally for a density matrix \rho is given by p(i) = Tr[P_i \rho P_i^{\dagger}]
Where P_i is the projection onto the eigenspace of O corresponding to \lambda_i
```

Can define more general measurements (non-projective)
 See next lecture!

- Quantum Computation and Quantum Information by Nielsen & Chuang: 2.1.7, 2.4
- Introduction to Quantum Cryptography by Thomas Vidick and Stephanie Wehner: chapter 2
- Quantum Information Theory by Mark M. Wilde: chapter 3