Simulation, Analysis, and Validation of Computational Models

- 5. Numerics of Non-Linear Systems -



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- Numerical integration
- Collision
- Percolation

- Asymptotic behavior of a nonlinear system, when (stable) fixed points, limit cycles and tori can be excluded in a deterministic, simple dynamical system.
- Sensitivity to initial conditions implies noise-like features, but short-term predictions are possible.
- Contains typically many unstable periodic trajectories.
- Occurs when the period of the trajectory of the system's state diverges to infinity.
- Occurs everywhere in nature, and in social and engineered systems.
- Mild chaos can add a bit of realism to a simulation.

H. Sayama (2015) Introduction to the Modeling and Analysis of Complex Systems. SUNY.

- Non-linear dynamics is predominant in all aspects nature, in society, and in complex technical systems.
- A major challenge in many fields of application of the theory is the identification of imminent qualitative changes ("tipping points").
- While linear systems are well understood, many problem in non-linear dynamics remain to be solved, to be formulated or to be discovered.
- Numerical simulation of non-linear systems are still improvable.

Numerical integration: Simulating non-linear dynamics

- Goal: If there is no analytical solution, we need to find a numerical solution x̂(t) for t ≥ t₀ for an ordinary differential equation ẋ = f(x), given x(t₀) = x̂₀(0) = x₀
- An algorithm that calculates \hat{x} can be evaluated by the
 - integral

$$\int_{t_0}^{t_1} \left(x\left(t\right) - \hat{x}\left(t\right) \right)^2 dt$$

- maximum of the difference $|x \hat{x}(t)|$
- maximal order of a polynomial that is perfectly be integrated.
- General idea: Find x at some points across the interval [t, t + Δt] using the given ODE and Taylor expansion. Resulting values are weighted and added to get an estimate for t + Δt.
- Integration formulas depend on the specific method and differ in number and weights and order of points along time axis.

Numerical integration

Euler method is the simplest integration formula

$$x(t + \Delta t) = x(t) + \dot{x}(t) \Delta t$$
 where $\dot{x}(t) = f(x(t))$

Perfect only for constant functions, but can be sufficient in other cases (for small Δt), Errors can be relatively large especially for exponential dynamics. Example: $\dot{x}(t) = x(t)$ starting at x(0) = 1 using $\Delta t = 1$ $x(1) = x(0) + x(0) \cdot 1 = 2$ $x(2) = x(1) + x(1) \cdot 1 = 2 + 2 \cdot 1 = 4$ $x(3) = x(2) + x(2) \cdot 1 = 4 + 4 \cdot 1 = 8 \Rightarrow x(n) = 2^{n}$ At n = 4 Euler: 16, exp(4) = 54.60: abs. error of 38.60 $\Delta t = 1$ is generally too large, but reducing step width to 0.1 still gives an abs. error of 9.34

Many better methods available.

Numerical integration: Predictor-corrector method

Result from a simple formula can be corrected by another method.

Combinations of any methods are possible, but for simplicity we start again with the Euler method for the Predictor step:

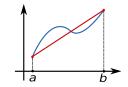
$$\tilde{x}(t + \Delta t) = x(t) + \dot{x}(t)\Delta t$$

Corrector step to interpolate the initial attempt with new result from the predictor using the differential equation $\tilde{\dot{x}} = f(\tilde{x}(t + \Delta t))$

$$\hat{x}(t + \Delta t) = x(t) + \frac{1}{2}\Delta t \left(\dot{x}(t) + f(\tilde{x}(t + \Delta t)) \right)$$

which is called trapezoidal rule.

Above example: $\tilde{x}(1) = 2$, $f(\tilde{x}) = 2 \Rightarrow \hat{x}(1) = 1 + \frac{1}{2} \cdot 1 \cdot (1+2) = 2.5 \ (\Delta t = 1)$ (compare to e = 2.71828).



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Numerical integration: Various classifications

- Explicit methods determine the state for a later time from the state of the system at present
- Implicit methods determine the solution by solving an equation involving the current state of the system and later ones.
- Direct methods: Use intermediate values to calculate next step
- Multi-step method: Calculate new values based on several previous values (needs to get started by simple method)

Numerical integration: Explicit methods

• E.g. RK4 (4th order Runge-Kutta method, explicit)

$$x(t + \Delta t) = x(t) + \frac{\Delta t}{6} \left(\underbrace{k_1 + 2k_2 + 2k_3 + k_4}_{\text{mixture of derivatives}} \right),$$

$$k_1 = f(x(t)),$$

$$k_2 = f\left(x\left(t + \frac{\Delta t}{2}\right) + \Delta t\frac{k_1}{2}\right),$$

$$k_3 = f\left(x\left(t + \frac{\Delta t}{2}\right) + \Delta t\frac{k_2}{2}\right),$$

$$k_4 = f\left(x(t + \Delta t) + \Delta t k_3\right).$$

• 4th order: The error of the result scales with $(\Delta t)^4 \times 4$ th derivative of the function f. I.e. it solves ODEs with polynomial f of order not greater than 3 precisely.

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Numerical integration: Implicit methods

• For highly non-linear ODEs, implicit methods may work better, i.e. moving forwards and backwards in time (like predictor corrector). Example: Adams-Moulton method.

$$\begin{aligned} x(t + \Delta t) &= x(t) + \frac{h}{720} \left(251f(x(t + \Delta t)) + 646f(x(t)) \right. \\ &\left. -264f(x(t - \Delta t)) + 106f(x(t - 2\Delta t)) \right. \\ &\left. -19f(x(t - 3\Delta t))) \right. \end{aligned}$$

- Need to get started (multi-step method), by calculating the first steps by more simple methods
- Solves ODEs with polynomial *f* of order not greater than 4 precisely (order 5).

Numerical integration: Points to remember

- Because analytical solutions are mostly unavailable, numerical integration is generally used.
- Trade-off between complexity of the integration formula and a small step width: Simply reducing step widths can give an idea that the system is problematic, but may be inefficient for solution.
 - Complex integration methods are usually available in numerical modelling
 - Practically, 4th order Runge-Kutta (RK4) is often sufficient.
 - Adaptive methods, step width control (not discussed here)
 - For complex or chaotic systems, implicit methods should be considered.
- Check: Is the energy conserved also in the simulated system?

- Collisions are essential in particle dynamics
- Should be avoided in robotics (i.e. occur at zero speed), but needed in robot simulator
- Assumptions: Elastic, point contact, no spin, no friction, no complex shapes, no deformation, no internal degrees of freedom.
- Trajectory at the collision is not smooth, because of instantaneous effect is assumed.
- Need to stop integration and handle collision separately.

Collisions in 1D

Two masses m_1 and m_2 on a line collide with velocities v_1 and v_2 and depart with velocities v'_1 and v'_2 (masses remain unchanged)

For an elastic collision, and total momentum

$$m_1v_1 + m_2v_2 = m_1v_1' + m_2v_2'$$

and total energy

$$\frac{1}{2}m_1v_1^2 + \frac{1}{2}m_2v_2^2 = \frac{1}{2}m_1v_1'^2 + \frac{1}{2}m_2v_2'^2$$

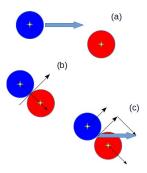
are conserved, although momentum and energy can be transferred from one object to the other one

I.e. weighted mean velocity + proportional slowdown or speedup (JFYI: https://www.youtube.com/watch?v=HEfHFsfGXjs)

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Collision in 2D

- (a) Relative motion v_1 of a moving ball against resting one $v_2 = 0$
- (b) Contact between the two convex objects defines coordinate system
- (c) After collision, original velocity splits according to contact coordinates



New v'_2 will be in direction from x_1 to x_2 (at collision): $x_2 - x_1$.

$$\begin{array}{l} \text{Project } \textbf{v}_1 \text{ on unit vector } \frac{x_2-x_1}{\|x_2-x_1\|} \colon \ \left\langle \textbf{v}_1, \frac{x_2-x_1}{\|x_2-x_1\|} \right\rangle \frac{x_2-x_1}{\|x_2-x_1\|} \text{ for } \textbf{v}_2 = 0 \\ \\ \text{ or } \ \left\langle \textbf{v}_1 - \textbf{v}_2, \frac{x_2-x_1}{\|x_2-x_1\|} \right\rangle \frac{x_2-x_1}{\|x_2-x_1\|} \text{ for any } \textbf{v}_2 \end{array}$$

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Collision in 2D

 $\begin{array}{l} \mathsf{v}_1' \text{ direction is obtained from the vector orthogonal to } (\mathsf{x}_2-\mathsf{x}_1).\\ \tilde{\mathsf{v}}_1 & \propto & \displaystyle \frac{\left< \mathsf{v}_1 - \mathsf{v}_2, (\mathsf{x}_2 - \mathsf{x}_1)^{\perp} \right>}{\|\mathsf{x}_2 - \mathsf{x}_1\|^2} \left(\mathsf{x}_2 - \mathsf{x}_1\right)^{\perp} & \left(\begin{array}{c} \mathsf{a} \\ \mathsf{b} \end{array} \right)^{\perp} = \left(\begin{array}{c} -\mathsf{b} \\ \mathsf{a} \end{array} \right)\\ \tilde{\mathsf{v}}_2 & \propto & \displaystyle \frac{\left< \mathsf{v}_1 - \mathsf{v}_2, \mathsf{x}_2 - \mathsf{x}_1 \right>}{\|\mathsf{x}_2 - \mathsf{x}_1\|^2} \left(\mathsf{x}_2 - \mathsf{x}_1\right) \end{array}$

The respective lengths are found as in the 1D case (see above)

$$\begin{array}{rcl} v_1' & = & v_1 + \frac{2m_2}{m_1 + m_2} \left(v_2 - v_1 \right) \\ v_2' & = & v_2 + \frac{2m_1}{m_1 + m_2} \left(v_1 - v_2 \right) \end{array}$$

So that we arrive at

$$\begin{aligned} \mathsf{v}_1' &= \mathsf{v}_1 + \frac{2m_2}{m_1 + m_2} \ \frac{\langle \mathsf{v}_2 - \mathsf{v}_1, (\mathsf{x}_2 - \mathsf{x}_1)^\perp \rangle}{\|\mathsf{x}_2 - \mathsf{x}_1\|^2} \ (\mathsf{x}_2 - \mathsf{x}_1)^\perp \\ \mathsf{v}_2' &= \mathsf{v}_2 + \frac{2m_1}{m_1 + m_2} \ \frac{\langle \mathsf{v}_1 - \mathsf{v}_2, \mathsf{x}_2 - \mathsf{x}_1 \rangle}{\|\mathsf{x}_2 - \mathsf{x}_1\|^2} \ (\mathsf{x}_2 - \mathsf{x}_1) \end{aligned}$$

Check conservation of momentum and energy!

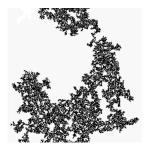
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- Sort objects lexicographically by coordinate
- Check first coordinate (then second coordinate etc.) for any close objects (or hierarchically for object parts) to fill active list
- $\bullet\,$ Calculate any position increments \dot{x} (see previous lecture) and check active list for penetration
- Revert increment of relevant object parts and invoke collision handler
- Calculate forces to other parts of the same object

- In a collision the (possibly) linear microscopic dynamics is abstracted to a nonlinear macroscopic effect
- Using heuristics and simple elementary shapes can be useful
- Realistic collision of deformable material is still subject of research
- More than two objects: Chaining, jamming, percolation

- Last lecture: Non-linear equations describing a qualitative change of a macroscopic system
- In phase transitions, order-parameters are used to describe typical phenomena
- Universality: Order parameter effects implied by typical equations (such as $\dot{x} = cx x^3$)
- Percolation as a simple example of a phase transition in microscopic description

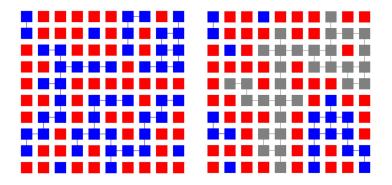
- Classical example: Mix an in increasing fraction of iron particle with sand, and check a whether the mixture conducts electricity
- Mathematical problem: Formation of a giant component



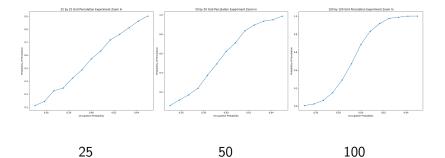
Other examples: Transport in porous media, underground CO_2 storage, drip-through coffee making, groundwater recharge, wear and tear, forest fires, neural avalanches, spread of diseases (later)

- $N \times N$ grid of cells
- Choose randomly whether a cell is conductive (probability p) or not (probability 1 - p)
- Use flooding algorithm to identify conductive cluster: All conductive cells with with a conductive cell next to them
- Repeat many times to determine probability of conductivity of the material as a whole
 - practically zero for small p
 - steeply increasing for a range of \boldsymbol{p}
 - practically one for large p
- Steepness of increase increases with N and becomes infinite for $N \rightarrow \infty$: Unique value p_{crit}
- Quantitative: Throughput increases for p > p_{crit}, similar to pitchfork bifurction (see last lecture)

Percolation in a 10 \times 10 square grid (top to bottom) near critical occupation probability p = 0.492



Percolation probability in an $N \times N$ square grid over occupation probability



- Practically interesting with a many applications
- Shows a common type of phase transition
- Distribution of the size of clusters at criticality decays as a power law with the same exponent for 2D lattices
- Percolation clusters are self-similar (*fractals*)
- When modelling disease spreading, percolation with be studied on a graph

- Noise
- System modelling
- Simulations